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Static Buckling and Free Vibration Analysis of Bi-Dimensional FG Metal Ceramic Porous Beam

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KEYWORDS

ABSTRACT

Porosity; Computational analysis; 2D-FGM; Composite materials; Buckling. This study presents an analytical solution for static buckling and free vibration analysis of bi-dimensional functionally graded (2D-FG) metal-ceramic porous beams. To achieve this goal, equations of motion for the beam are derived by using Hamilton's principle and then the derived equations were solved in the framework of Galerkin's well-known analytical method for solution of equations. The material properties of the beam are variable along with thickness and length according to the power-law function. During the fabrication of functionally graded materials (FGMs), porosities may occur due to technical problems causing micro-voids to appear. Detailed mathematical derivations are presented and numerical investigations are performed, while emphasis is placed on investigating the effect of various parameters such as FG power indexes along both directions of thickness and length, porosity, and slenderness ratios (L/h), on the non-dimensional frequency and static buckling of the beam based on new higher deformation beam theory. The accuracy of the proposed model is validated based on comparisons of the results with the accepted studies. According to the result in both buckling and vibration analysis, the presented modified transverse shear stress along the thickness has shown closer consequences in comparison with TBT.

1. Introduction

In recent years, a new class of composite materials formed that contained two or more phases which are known as functionally graded materials (FGMs). FGMs are made of two or more whose characteristics substances vary continuously from one direction to another. The performances and advantages of FGMs have led to their increasing use in many sectors including hi-tech industries, aerospace, naval, automotive, and civil structures[1]. An important application of FGM in structural components is in beams and beam structures. To support the growing use of FGM beams, the study of their vibration behavior becomes necessary. In recent years, the analysis of the free vibration of FGM beams has been performed using different analytical and numerical methods[2]. buckling and free vibration of a curved Timoshenko FG microbeam is studied by Rahmani et al. [3] based on strain gradient theory (SGT) theory. Compared to

conventional materials, porous materials possess both remarkable design flexibility and significant applicability across diverse fields. Consequently, numerous researchers have dedicated their efforts to exploring various porous structures. For instance, Ghasemi and Meskini [4] conducted an investigation into the free vibration analysis of porous laminated rotating circular cylindrical shells. Several studies have been carried out on the behaviors of functionally graded material in recent years [5-10]. Atashipour et al. [11] developed a closed-form 2D elasticity solution for stresses and displacements of a curved FG beam subjected to a shear force. Among the earliest investigators in this field, Masjedi et al. [12], conducted a large deflection of functionally graded porous beams based on a geometrically exact theory. The buckling behavior of engineering structures made of advanced materials has been investigated by many researchers[13-18]. Buckling of Sandwich Structures with Metamaterials Core Integrated

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by Graphene Nanoplatelets Reinforced Polymer Composite was analyzed by Shabani and Khorshidi[19]. The analytical solution presented by Dym and Williams [20] can be used to estimate the buckling behavior of curved beams. The nondeterministic vibration frequencies and mode shapes of FG porous beams were investigated by Gao et al. [21] using a hybrid Chebyshev surrogate model with a discrete singular convolution method. Large amplitude nonlinear vibration analysis of functionally graded porous Timoshenko beams is examined by Ebrahimi and Zia [22]. Chen et al. [23] investigated free vibration and transient analyses of functionally graded piezoelectric materials curved Timoshenko beams in their work. Mahmoodabadi et al. [24] used a multi-objective optimization process based on the genetic algorithm to study damping the vibrations of a piezo-actuating composite beam. The Peridynamics analytical solutions for the buckling load of beams with various boundary conditions were introduced by Yang et al. [25]. Ly et al. [26] used different hybrid ML models to predict the critical buckling load of I-shaped cellular steel beams with circular openings.

In this study, the effects of variable L/h, porosity coefficient, and gradient indexes in two directions on natural frequency and critical buckling load of 2D-FG porous beams have been investigated. Young's modulus, mass density, and Poisson's ratio of the beam are assumed to vary along the thickness and length of the beam according to power-laws form and also, they are influenced by the porosity. In this study to achieve better accuracy, modified transverse shear stress has been presented and compared with other theories. In the following, the higherorder governing equations are derived by using Hamilton's principle. In the end, the Galerkin method is employed to solve them, and the nondimensional frequencies and critical buckling loads are obtained.

2. Problem and Formulation

2.1. Two Dimensionally Functionally Graded Porous Beams

Consider a beam as shown in Fig. 1. with length L, width b, and thickness h, with Cartesian coordinate system $O(x \ y \ z)$ which the origin of coordinate system O is chosen at the left of the beam. The Mechanical properties of the beam, such as Young's modulus E(x, z), shear modulus G(x, z), Poisson's modulus v(x, z) and mass density $\rho(x, z)$ with the material properties can be varying along length and thickness like Fig. 1.



The effective material properties(P) can be expressed by using the rule of mixture:

$$P(x,z) = P_{c}R_{c}(x,z) + P_{m}R_{m}(x,z)$$
(1)

$$R_{c}(x,z) + R_{m}(x,z) = 1$$
(2)

where P_c and P_m are epitomes of mechanical properties and R_c and R_m are volume fractions of ceramic and metal. Substitution of Eq. (2) into Eq. (1), the material properties of the 2D-FG beam are obtained as [27]:

$$P(x,z) = (P_c - P_m)R_c + P_m$$
(3)

As shown in Fig. 2, it is assumed that the porosities in the x and z directions are distributed uniformly. The effective material properties of the even porosity are defined as:

$$P(x,z) = (P_c - P_m)R_c + P_m - \left(\frac{\eta}{2}(P_c + P_m)\right)$$
(4)

where η is a porous parameter or porosity volume fraction.

Fig. 2. Porosity of beam.

The assumed beam is graded from metal at the lower left corner edge to ceramic at the top right corner edge (Fig. 1). The volume fraction of ceramic material is given by:

$$R_{c}(x,z) = \left(\frac{x}{L}\right)^{k_{x}} \left(\frac{z}{h} + \frac{1}{2}\right)^{k_{z}}$$
(5)

The k_x and k_z are the power-law of the beam which ascertain properties in length and thickness direction. The mechanical properties of the 2D-FG porous beam can be written as [28]:

$$E(x,z) = (E_c - E_m) \left(\frac{x}{L}\right)^{k_x} \left(\frac{z}{h} + \frac{1}{2}\right)^{k_z} + E_m - \left(\frac{\eta}{2}(E_c + E_m)\right)$$
(6a)

$$\rho(x,z) = (\rho_c - \rho_m) \left(\frac{x}{L}\right)^{k_x} \left(\frac{z}{h} + \frac{1}{2}\right)^{k_z} + \rho_m - \left(\frac{\eta}{2}(\rho_c + \rho_m)\right)$$
(6b)

$$\upsilon(x,z) = (\upsilon_c - \upsilon_m) \left(\frac{x}{L}\right)^{k_x} \left(\frac{z}{h} + \frac{1}{2}\right)^{k_z} + \upsilon_m - \left(\frac{\eta}{2}(\upsilon_c + \upsilon_m)\right)$$
(6c)

2.2. Mathematical Modeling

The displacement filed the present shear deformation beam theories, are given as [29]:

$$u_{x}(x,z) = u(x) - g(z)\frac{\partial w(x)}{\partial x} + f(z)\varphi(x)$$
 (7a)

$$u_{z}(x,z) = w(x)$$
(7b)

where u(x) and w(x) respectively represent the axial and transverse displacements for the midaccess, and φ is the rotation of the cross-sections. Assuming infinitesimal deformations, straindisplacement relations are[30]:

$$\varepsilon_{xx} = \frac{\partial u(x)}{\partial x} - g(z)\frac{\partial^2 w(x)}{\partial x^2} + f(z)\frac{\partial \varphi(x)}{\partial x}$$
(8a)

$$\gamma_{xz} = -\frac{\partial g(z)}{\partial z} \frac{\partial w(x)}{\partial x} + \frac{\partial f(z)}{\partial z} \varphi(x) + \frac{\partial w(x)}{\partial x}$$
(8b)

The stress-strain relations by using Hook's law can be defined as follows:

$$\sigma_{\rm XX} = E(x, z) \varepsilon_{\rm XX} \tag{9a}$$

$$\tau_{\rm XZ} = k_s G(x, z) \gamma_{\rm XZ} \tag{9b}$$

E and G represent the modulus of elasticity and shear moduli, respectively, while ks denotes the shear correction factor. Where the shear modulus (G) is [31]:

$$G(x,z) = \frac{E(x,z)}{2(1+\nu(x,z))}$$
(10)

Hamilton's principle is employed to extract equations of motion [32, 33].

$$\int_{0}^{t} \left(\delta U - \delta T \right) dt = 0 \tag{11}$$

where U and T are the strain and kinetic energy, respectively of the beam. δ is the variation operator. The strain energy of the beam(U) is calculated as follows [34, 35]:

$$U = \frac{1}{2} \int_{V} \sigma_{ij} \varepsilon_{ij} dV = \frac{1}{2} \int_{V} (\sigma_{xx} \varepsilon_{xx} + \tau_{xx} \gamma_{xx}) dV$$
(12)

Finally, variation of strain energy with respect to u(x), w(x) and $\phi(x)$ is shown as.

$$\delta U = \frac{1}{2} \int_{A} \left[-\delta u \, \frac{\partial A_{xx}}{\partial x} - \delta w \, (\frac{\partial^2 B_{xx}}{\partial x^2} - \frac{\partial S_{xz}}{\partial x^2} + \frac{\partial A_{xz}}{\partial x}) - \delta \varphi (\frac{\partial D_{xx}}{\partial x} - T_{xz}) \right] dA$$
(13)

where A_{XX} , A_{XZ} , B_{XX} , D_{XX} , S_{XZ} , and T_{XZ} are defined by:

$$(A_{xx}, A_{xz}) = -\int_{-h/2}^{h/2} (\sigma_{xx}, \tau_{xz}) dz$$
 (14a)

$$(B_{xx}, D_{xx}) = \int_{-h/2}^{h/2} (g(z), f(z)) \sigma_{xx} dz$$
 (14b)

$$(\mathbf{S}_{\mathbf{X}\mathbf{Z}}, \mathbf{T}_{\mathbf{X}\mathbf{Z}}) = \int_{-h/2}^{h/2} (\frac{\partial \mathbf{g}(\mathbf{z})}{\partial \mathbf{z}}, \frac{\partial \mathbf{f}(\mathbf{z})}{\partial \mathbf{z}}) \tau_{\mathbf{X}\mathbf{Z}} d\mathbf{z} \qquad (14c)$$

The kinetic energy is obtained as [36]:

$$T = \frac{1}{2} \int_{V} \rho \left[\dot{u}_{x}^{2} + \dot{u}_{z}^{2} \right] dV$$
 (15)

The inertia coefficients appearing in Eq. 19 can be defined as:

$$(I_1, I_2, I_3, I_4, I_5, I_6) = -\int_{-h/2}^{h/2} \rho(1, g(z)),$$

$$f(z), g(z) f(z), g(z)^2, f(z)^2) dz$$
(16)

Finally, the total variation of kinetic energy associated with the sandwich beam in the integral form as:

$$\delta T = \frac{1}{2} \int_{A} \left[\left(-I_{1} \ddot{u} + I_{2} \frac{\partial \ddot{w}}{\partial x} - I_{3} \ddot{\varphi} \right) \delta u + \left(-I_{3} \ddot{u} + I_{4} \frac{\partial \ddot{w}}{\partial x} - I_{6} \ddot{\varphi} \right) \delta \varphi - \left(I_{2} \frac{\partial \ddot{u}}{\partial x} - I_{5} \frac{\partial^{2} \ddot{w}}{\partial x^{2}} + I_{4} \frac{\partial \ddot{\varphi}}{\partial x} - I_{1} \ddot{w} \right) \delta w \right] dA$$

$$(17)$$

By substituting strain energy Eq. (13) and kinetic energy Eq. (17) into Hamilton's principal Eq. (11), equations of motion may be expressed as Eqs. (18a) – (18c).

$$\delta u : \frac{\partial A_{xx}}{\partial x} = I_1 \ddot{u} - I_2 \frac{\partial \ddot{w}}{\partial x} + I_3 \ddot{\varphi}$$
(18a)

$$\delta w : \frac{\partial^2 B_{xx}}{\partial x^2} - \frac{\partial S_{xz}}{\partial x} + \frac{\partial A_{xz}}{\partial x} = I_2 \frac{\partial \ddot{u}}{\partial x}$$

$$- I_5 \frac{\partial^2 \ddot{w}}{\partial x^2} + I_4 \frac{\partial \ddot{\varphi}}{\partial x} + I_l \ddot{w}$$
 (18b)

$$\delta \varphi : \frac{\partial \mathbf{D}_{\mathbf{x}\mathbf{x}}}{\partial x} - \mathbf{T}_{\mathbf{x}\mathbf{z}} = \mathbf{I}_3 \ddot{u} - \mathbf{I}_4 \frac{\partial \ddot{w}}{\partial x} + \mathbf{I}_6 \ddot{\varphi}$$
(18c)

2.3. Analytical Solution

To obtain the theoretical solution, the Galerkin method is considered. According to this method, the displacements functions u(x, t), w(x, t) and $\varphi(x, t)$ are assumed as follows[37]:

$$u(x,t) = \sum_{m=1}^{K} \left[Y_m \bar{u}_m \right] e^{i \, \omega t}$$
(19a)

$$w(x,t) = \sum_{n=1}^{K} [\Psi_n \overline{w_n}] e^{i\omega t}$$
(19b)

$$\varphi(x,t) = \sum_{j=1}^{K} \left[\Phi_{j} \,\overline{\varphi}_{j} \right] e^{i \, \omega t} \tag{19c}$$

where \bar{u}_m , \bar{w}_n and $\bar{\varphi}_j$ are unknown coefficients that will be determined. i = $\sqrt{-1}$, K denotes the order of series and ω is the natural frequency. $Y_m(x)$, $\Psi_n(x)$ and $\Phi_j(x)$ are the admissible functions that satisfy the fully clamped boundary conditions.

$$Y_m(x) = (L - x)^{q_0} x^{(m+p_0)-1}$$
(20a)

$$\Psi_n(x) = (L - x)^{q_0} x^{(n+p_0)-1}$$
(20b)

$$\Phi_{j}(x) = (L-x)^{q_{0}} x^{(j+p_{0})-1}$$
(20c)

Free vibration analysis of the bi-dimensional functionally graded sandwich beam can be computed from Eq. (21a)[38]. Also, the amount of critical buckling loads will be calculated by Eq. (21b):

$$\left\{ \begin{bmatrix} K \end{bmatrix} - \omega^2 \begin{bmatrix} M \end{bmatrix} \right\} \left\{ \lambda \right\} = 0$$
 (21a)

$$\left\{ \begin{bmatrix} K \end{bmatrix} - N_{cr} \begin{bmatrix} K_g \end{bmatrix} \right\} \left\{ \lambda \right\} = 0$$
 (21b)

where M is the global mass matrix, K is the stiffness matrix and unknown coefficients of Eq. (21) is $\{\lambda\}$.

3. Numerical Results and Discussion

In this section, the free vibration of the 2D-FG porous beam with respect to porosity coefficients (η) is studied. Functionally graded material composed of a mixture of alumina and steel as ceramic and metal respectively with the material. Their properties are given in Table 1. The influence of different slenderness ratios, L/h = 5, 10, 15, and 20 on the non-dimensional natural frequency are investigated. The shear correction factor for TBT theory is considered as ks = 5/6 [39].

Table 1. Properties of materials[37]

	=		
Material	Elasticity module (E)	Mass density (ρ)	Poisson's ratio (υ)
Alumina	380	3800	0.23
Steel	70	2700	0.23

The dimensionless fundamental frequency and critical buckling load are defined as Eq. 22.

$$\overline{\omega} = \omega^2 \sqrt{\frac{\rho_c A}{E_c I}}$$
(22a)

$$\overline{N}_{cr} = N_{cr} \frac{12L^2}{E_c h^3}$$
(22b)

where I represent the moment of inertia and A represents the area of the cross-section of the beam.

$$I = \frac{A \cdot h^2}{12} \tag{23}$$

In this research, the solution is calculated for different the power-law index between 0 and 10 moreover, porosity coefficients are taken as $\eta = 0, 0.1, \text{ and } 0.2$. The total thickness of the beam (h) is constant and it is 0.1 m. The function indexes (p0 and q0), are taken as 2 for satisfying the clamped-clamped boundary condition. Besides, g(z) and f(z) are related parameters to the beam's theory, as they take as 0 and z for TBT, and z and $z\left(\frac{4}{h}-\frac{16z^2}{3h^3}\right)$ for NHOBT [40], respectively. Validation of our formulation and the results are obtained and compared with the results of Ref. [41], [42], and [43]. In Table 2, the first frequency of FG porous less beam with L/h = 5 and $k_x = 0$ are calculated and they will be compared with references. To verify our new theory, the natural frequency of the FG beam with a greater slender ratio in Table 3 is calculated.

In addition, another comparison with ref [41] and [43] is presented to show the accuracy of the new higher shear deformation beam theory to calculate the critical buckling load. The mechanical properties of the pure beam have changed along the z direction and the ratio of beam's length to height is 10. The results are written in Table 4.

The amounts of Ncr of the FG beam with NHOBT and L/h = 5 are reported in Table 5. By checking both tables, it is clear that there is accuracy and agreement between the references and the present study. Thus, the results of natural frequencies and critical buckling loads will be published under effective parameters such as L/h, porosity (η), and functionally graded power indexes (k_x and k_z).

The results of non-dimensional frequencies of a two-dimensional functionally graded porous beam with several η , k_x , and k_z by using two theories are written in Table 6. The amount of ϖ has been reduced by approximately 37% by changing k_z from 0 to 8. On the other hand, the variation of k_x has affected ϖ by around 34%, indicating that k_z has a greater impact on the results. Additionally, by changing η from 0 to 0.2, the amount of ϖ has increased by approximately 3%. The amount of Ncr under changing k_x , k_z , and η with TBT and NHOBT is calculated and the consequences are written in Table 7. By changing k_z and k_x between 0 to 8, the amount of Ncr has reduced by about 76 and 69 % respectively. The factor with the greatest impact on the results is k_z , with approximately 7% more effect. Furthermore, the Ncr decreased by about 11% with increasing η from 0 to 0.2.

Table 2. Comparison of the dimensionless natural frequencies $(l/h = 5 \text{ and } \eta = 0)$								
			kz =	0	kz = 1	kz = 2	kz = 5	kz = 10
Present (NH	HOBT)		10.2	285	8.0969	7.3290	6.6448	6.3042
Present (TB			10.2579		8.1098	7.3750	6.7128	6.4759
Thuc et al. H		1	10.0	678	7.9522	7.1801	6.4961	6.1662
Thuc et al. F			9.99	84	7.9015	7.1901	6.6447	6.3161
Simsek HOE			10.0	705	7.9503	7.1767	6.4935	6.1652
Simsek FOB			10.0		7.9253	7.2113	6.6676	6.3406
		Table 3.	Comparison of	the dimension	lless natural frequ	uencies (l/h = 20	and η = 0)	
			kz =	0	kz = 1	kz = 2	kz = 5	kz = 10
Present (NH	HOBT)		12.2	418	9.4918	8.6778	8.2159	7.9342
Thuc et al. H		1	12.2		9.4328	8.5994	8.1460	7.8862
Thuc et al. F			12.2		9.4311	8.6047	8.1698	7.9115
Simsek HOBT			12.22		9.4315	8.5975	8.1446	7.8858
Simsek FOBT			12.22	235	9.4314	8.6040	8.1699	7.9128
Table 4. Comparison of the results of critical buckling load ($l/h = 10$ and $\eta = 0$)								
			kz =	0	kz = 1	kz = 2	kz = 5	kz = 10
Present (NH	HOBT)		194.	.4026	98.8053	76.8331	62.9629	56.4158
Li and Batra	a [43]		195.	.3400	98.7490	76.9800	64.0960	57.7080
Thuc et al. F	FOBT [41]		195.	.3730	98.7923	77.0261	64.1324	57.7329
Thuc et al. H	HOBT [41]]	195.	.3610	98.7868	76.6677	62.9786	56.5971
		Table	5. Comparison	of the results o	of critical bucklin	g load (l/h = 5 an	ıd η = 0)	
			kz =	0	kz = 1	kz = 2	kz = 5	kz = 10
Present (NH	HOBT)		152.	.1990	79.8063	61.2369	47.116	41.1128
Li and Batra	a [43]		154.	.3500	80.4980	62.6140	50.3840	44.2670
Thuc et al. F	FOBT [41]		154.	154.4150		62.6616	50.4207	44.2946
Thuc et al. HOBT [41]		154.	.5500	80.6087	61.7925	47.7562	41.8042	
Table 6. First non-dimensional frequency (ϖ) of 2D-FG porous beam for various kx, kz, and η (L/h = 10). kx Theory η kz						10).		
	•		0	1	2	4	6	8
		0	21.6960	17.0553	15.6524	14.5606	14.0223	13.6403
		1	16.8500	14.3819	13.6517	13.0863	12.7849	12.5764
TBT 0	2	15.3637	13.6340	13.0914	12.6367	12.4165	12.2903	
	4	14.6350	13.2308	12.7760	12.3834	12.1918	12.0892	
	6	14.3520	13.0449	12.6168	12.2521	12.0960	11.9658	
		8	14.1373	12.8873	12.4975	12.1511	12.0043	11.9063
	0	21/1618	16/6429	15/3731	14/2928	13/7558	13/3948	
		1	16/5051	14/0453	13/3743	12/8050	12/5325	12/3574
		2	15/0571	13/3139	12/8158	12/3854	12/1817	12/0524
NHOBT	0	4	14/2740	12/8894	12/4747	12/1139	11/9551	11/8567
		6	13/9665	12/6941	12/3101	11/9907	11/8451	11/7547
	8	13/7419	12/5464	12/1924	11/8892	11/7570	11/6799	

		0	21.9998	16.7114	15.0772	13.6923	13.0121	12.5896
TBT 0.1	1	16.5388	13.6219	12.7145	12.0111	11.6668	11.4278	
	2	14.6162	12.6623	12.0187	11.4848	11.2217	11.0879	
IDI	0.1	4	13.6436	12.1432	11.6160	11.1728	10.9759	10.8546
		6	13.3410	11.9406	11.4468	11.0374	10.8555	10.7441
		8	13.1444	11.7965	11.3355	10.9355	10.7829	10.6676
		0	22/3579	16/9650	15/3647	14/0096	13/3560	12/9310
		1	16/7512	13/8084	12/9738	12/2489	11/9219	11/7106
NUODT	0.1	2	14/8641	12/8713	12/2623	11/7220	11/4849	11/3378
NHOBT	0.1	4	13/9234	12/3669	11/8677	11/4237	11/2257	11/1112
		6	13/6599	12/1858	11/7091	11/2917	11/1181	11/0058
		8	13/4683	12/0397	11/5834	11/1987	11/0331	10/9345
		0	22.3718	16.2578	13.9010	12.1407	11.3053	10.7474
		1	16.0382	12.4432	11.2190	10.2780	9.8109	9.5400
трт	0.2	2	13.4114	11.1096	10.2812	9.5863	9.2764	9.0820
TBT	0.2	4	11.8584	10.3262	9.7099	9.1750	8.9410	8.8175
		6	11.4574	10.0734	9.5067	9.0297	8.8242	8.6936
		8	11.3014	9.9231	9.3867	8.9382	8.7464	8.6273
		0	22/7021	16/3349	14/2677	12/4383	11/6037	11/0778
NHOBT 0.2		1	16/2534	12/5477	11/4541	10/4811	10/0477	9/7799
	0.2	2	13/6285	11/2567	10/4767	9/7967	9/4873	9/3149
	0.2	4	12/0916	10/4981	9/9149	9/3829	9/1574	9/0258
		6	11/7261	10/2714	9/7284	9/2452	9/0443	8/9243
		8	11/5545	10/1302	9/6076	9/1525	8/9648	8/8529

Table 7. Critical buckling	load (Ncr) of 2D-FG por	rous beam for va	arious kx. kz.	and n (L/h = 10).
Tuble 7. Gritteat Ducking	s iouu (nici	101 LD 1 U poi		1110u3 KA, KL,	$ana \eta (b) n = 10$

Theory η	n	1	kx					
	kz	0	1	2	4	6	8	
TBT 0		0	201.3001	101.2443	73/3782	56/8435	51/2608	48/3640
		1	101.8384	66.6279	55/4769	47/7935	44/9744	43/4770
	0	2	79.3865	58.3191	50/8079	45/2804	43/1927	42/0855
	4	68.2815	53.5313	47/9427	43/6789	42/0411	41/1664	
		6	64.3308	51.4026	46/5823	42/8864	41/4587	40/6963
		8	61.6250	49.8626	45/5830	42/2954	41/0188	40/3319
		0	195/3779	97/2664	71/2192	55/4709	50/1329	47/3546
		1	99/2310	64/5146	53/9762	46/6343	43/9498	42/5304
NUODT	0	2	77/1877	56/4267	49/3866	44/1401	42/1719	41/1327
NHOBT	0	4	65/7422	51/5953	46/4741	42/5032	40/9942	40/1957
		6	61/5646	49/4644	45/1055	41/7020	40/4027	39/7119
		8	58/8715	47/9899	44/1322	41/1171	39/9625	39/3461
		0	189/7017	86/5973	58/2271	42/5622	37/4739	34/8816
		1	88/0537	52/9688	41/9996	34/7625	32/1893	30/8440
TDT	0.1	2	64/0969	44/5446	37/5045	32/4321	30/5622	29/5826
TBT	0.1	4	52/5693	39/8829	34/8210	30/9839	29/5379	28/7694
		6	49/1145	38/0454	33/6697	30/3341	29/0618	28/3879
		8	46/9595	36/7694	32/8503	29/8570	28/7108	28/1008
		0	184/1132	82/7353	56/4400	41/5326	36/6694	34/1715
		1	85/8718	51/1773	40/8558	33/9401	31/4780	30/1892
NULODE	0.1	2	62/4375	43/0432	36/4529	31/6364	29/8622	28/9356
NHOBT	0.1	4	50/6270	38/3560	33/7335	30/1575	28/8153	28/1133
		6	46/8667	36/4948	32/5662	29/4938	28/3337	27/7230
		8	44/6357	35/2597	31/7613	29/0213	27/9825	27/4327
		0	178/0717	70/4071	40/9069	26/8716	22/6228	20/5068
		1	73/6780	38/2597	27/4892	20/9948	18/8072	17/6832
	0.0	2	47/4333	29/7059	23/3508	18/9993	17/4592	16/6723
TBT	0.2	4	34/6064	25/0341	20/8595	17/7380	16/5918	16/0017
		6	31/5006	23/5343	19/9695	17/2607	16/2562	15/7359
		8	30/0213	22/6233	19/3964	16/9431	16/0272	15/5484
		0	172/8112	66/3052	39/5456	26/1841	22/1330	20/0973
		1	71/9500	36/7506	26/6959	20/4944	18/4025	17/3281
NULODE	0.0	2	46/3571	28/6013	22/6763	18/5411	17/0741	16/3237
NHOBT 0.	0.2	4	33/4452	23/9716	20/1835	17/2666	16/1973	15/6521
		6	30/0084	22/4247	19/2647	16/7747	15/8535	15/3803
		8	28/3175	21/5147	18/6904	16/4552	15/6238	15/1932

The beam is made of two-dimensional functionally graded material; thus, it has been influenced by FG power indexes in x (k_x) and z (k_z) directions. The critical buckling loads (Ncr) and the natural frequencies (ϖ) of the beam are related to the k_x and k_z. In Figure 3a, the relation between ϖ and FG power indexes in a porous less beam is shown, the ratio of length to the height of the beam is considered as 10. In Figure 3b, the effect of FG power indexes on Ncr of the porous beam is under new theory (NHOBT) with $\eta = 0.2$ and L/h = 10 illustrated. Generally, the amount of Ncr and ϖ have reduced with increasing the FG power indexes. However, Ncr has been affected more than ϖ by changing k_x and k_z . In the high amount of FG power indexes the variations of Ncr have been moved to near 0. On the other hand, the ϖ has been diminished continuously by reducing the k_x and k_z. Power indexes affect how material properties vary along the length and thickness directions, which influence their stiffness and mass distribution. A higher FG power index reduces the effective stiffness and increases the effective mass of a beam.



Fig. 3a. Non-dimensional frequency of 2D-FG porous beam for various k_z and k_x (TBT, L/h = 10 and η = 0)



Fig. 3b. Critical buckling load of 2D-FG porous beam for various k_z and k_x (NHOBT, L/h = 10 and η = 0.2)

The impact of various ratios of the beam's spin to its height on ϖ and Ncr is shown in Figure 4. A beam with n = 0.1 and k_x = 1 is used to analyze the three amounts of L/h = 5, 10, and 20 on the ϖ in the first part of the Figure. A direct connection is seen between the amount of L/h and the ϖ . Additionally, the same study on Ncr has been conducted under unique conditions using a new theory. The results of this study are plotted in Figure 4b and the amount of Ncr has been rose by growing the L/h from 5 to 20. Also, the effect of L/h = 5 on the Ncr and ϖ is higher than 10 and 20. Reducing the strain energy will decrease the bending moment of the beam, so when the L/h is going to be diminished, the amount of Ncr and ϖ will be reduced. A higher slenderness ratio means a longer and thinner beam, which is more flexible and less stiff than a shorter and thicker beam.



Fig. 4a. Comparison of non-dimensional frequency of 2D-FG porous beam for various slenderness ratios (L/h) (η = 0.1 and k_x = 1)



Fig. 4b. Comparison of Critical buckling load of 2D-FG porous beam for various height to span ratios (L/h) (NHOBT, $\eta = 0.1$ and $k_x = 1$)

Another important and effective parameter is porosity. Figure 5 shows the effect of $\eta = 0, 0.1$, and 0.2 on the Ncr and ϖ . When the amount of η is 0, the beam does not have any porosity. In the constant k_z and k_z , the ϖ has shown an inverse direction with η so as to by increasing the η , ϖ has fallen. The same relation is seen between η and Ncr, however, the effect on Ncr is bolder than on ϖ . By reducing the stiffness of the beam, the amount of ϖ and Ncr will be influenced and decreased. Porosity affects the material density and strength along the thickness direction of beams, which in turn affects their stiffness and mass. Higher porosity means more voids or holes in the material, which reduces its weight and resistance to deformation. A higher porosity reduces the effective stiffness and increases the effective mass of a beam under both axial and transverse loading, which lowers its critical buckling load and natural frequency.



Fig. 5a. Effect of various porosity(η) on non-dimensional frequency of 2D-FG porous beam (L/h = 10 and kx = 1)



Fig. 5b. Effect of various porosity(η) on Critical buckling load of 2D-FG porous beam (NHOBT, L/h = 10 and kx = 8)

4. Conclusions

In this research paper, we have investigated the impact of various parameters, including porosity (η), slenderness ratios (L/h), and power-law indexes (kx and kz) in both axial and thickness directions, on the natural frequencies and critical buckling load of 2D-FG porous beams. We have utilized a new higher-order beam theory to analyze these effects.

Our findings demonstrate that the power-law indexes (kx and kz) play a significant role in determining the natural frequencies. Increasing the FG power indexes leads to a decrease in the values of the natural frequencies (Ncr) and the critical buckling load (ϖ) . This can be explained by the change in material properties along the beam's axial and thickness directions. The power-law indexes alter the distribution of the constituent materials, affecting their stiffness and density. Higher power-law indexes result in a more significant variation of material properties, leading to decreased natural frequencies and critical buckling loads. In the case of high powerlaw indexes, the variations in Ncr do not exhibit significant changes. Additionally, when considering a constant value for kx and kz, we observe an inverse relationship between the porosity (η) and ϖ . As η increases, ϖ decreases, and the same trend is observed between η and Ncr. This behavior can be attributed to the presence of voids or pores within the material. The presence of voids weakens the overall

structural integrity and reduces the effective stiffness of the beam. As a result, the natural frequencies decrease, and the critical buckling load is diminished. Furthermore, frequencies and critical buckling loads are more sensitive to porosity in high power-law indexes. A direct connection is seen between the amount of L/h and the ϖ . The amount of Ncr has been rose by growing the L/h from 5 to 20. Also, the effect of L/h = 5 on the Ncr and ϖ is higher than 10 and 20.

These findings provide valuable insights into the behavior of bi-dimensional FG metal ceramic porous beams and highlight the importance of considering porosity, slenderness ratios, and power-law indexes in their design and analysis. Further research in this area can contribute to the development of more efficient and reliable structures with improved performance characteristics.

Conflicts of Interest

The author declares that there is no conflict of interest regarding the publication of this manuscript. In addition, the authors have entirely observed the ethical issues, including plagiarism, informed consent, misconduct, data fabrication and/or falsification, double publication and/or submission, and redundancy.

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